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From instantons to quilts with seam degenerations

David L. Duncan

Michigan State University

2014 CMS Winter Meeting

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Σ ; closed, connected, oriented surface with Riemannian metric

$P \rightarrow \Sigma$; principal $SO(3)$ -bundle

Assume P is non-trivial. Then there are no reducible flat connections.

(More generally, this entire discussion carries over with $PU(r)$ -bundles, $r \geq 2$, provided the degree is coprime to r .)

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$$M(\Sigma) = \frac{\{\text{flat connections}\}}{\text{gauge}}$$

Theorem (Dostoglou-Salamon (1993))

$$HF_{\text{inst}}(S^1 \times \Sigma) \cong HF_{\text{symp}}(S^1 \times \Sigma)$$

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$$\mathbb{R} \times S^1 \rightarrow M(\Sigma),$$

where

$$M(\Sigma) = \frac{\{\text{flat connections}\}}{\text{gauge}}$$

The proof reduces to defining a bijection

$$\begin{aligned} & \{(\text{isolated}) \text{ holo. maps } \mathbb{R} \times S^1 \rightarrow M(\Sigma)\} \\ & \longrightarrow \{(\text{isolated}) \text{ instantons on } \mathbb{R} \times S^1 \times \Sigma\} / \text{gauge} \end{aligned}$$

for suitable auxiliary data (metric, perturbation).

Dostoglou-Salamon define this map using an implicit function theorem.

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Their theorem continues to hold if $\mathbb{R} \times S^1 \times \Sigma$ is replaced by $\mathbb{R} \times (\text{mapping torus})$.

The perspective of this talk is to view their theorem as a statement about 4-manifolds.

Question: Can we extend this to more general manifolds?

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Summary of results

Replace $\mathbb{R} \times S^1 \times \Sigma$ with an oriented, connected 4-manifold Z with cylindrical ends (allowing the possibility of no ends).

Replace $\mathbb{R} \times S^1$ with an oriented surface S having the same number of cylindrical ends.

Assume there is a map $F : Z \rightarrow S$ that preserves the ends and has connected, non-empty fibers.

Assume there is an $SO(3)$ -bundle $R \rightarrow Z$ that restricts to the non-trivial bundle on some fiber of F .

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Assume for now that F is a fiber bundle.

The more general case will be discussed briefly at the end.

For $x \in S$, let $\Sigma_x = F^{-1}(x)$ be the fiber over x .

Let $M(Z)$ be the fiber bundle over S with fiber over $x \in S$ given by $M(\Sigma_x)$.

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Theorem (D. (2014))

There is an injection

$$\mathcal{M} : \{(isolated) \text{ holo. sections } S \rightarrow M(Z)\} \\ \longrightarrow \{(isolated) \text{ instantons on } Z\} / gauge$$

for suitable auxiliary data (metric, perturbation).

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for suitable auxiliary data (metric, perturbation).

Either use the Dostoglou-Salamon implicit function theorem, or use the Yang-Mills heat flow. (The latter generalizes to when F is not a fiber bundle.)

Summary of results

Theorem (D. (2013))

There is a map

$$\begin{aligned} & \{(isolated) \text{ holo. sections } S \rightarrow M(Z)\} \\ & \longleftarrow \{(isolated) \text{ instantons on } Z\} / gauge : \mathcal{N} \end{aligned}$$

that is well-defined for the same auxiliary data (metric, perturbation) as \mathcal{M} .

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Theorem (D.-McNamara (2014))

The map \mathcal{N} is an injection.

Summary of results

Corollary

The relative Donaldson invariant of (Z, R, F) agrees with the relative symplectic invariant of (Z, R, F) .

Generalizes Dostoglou-Salamon's theorem, and gives an alternate proof.

'Relative' refers to the fact that the invariants take values in the Floer homology of the ends.

Should extend to polynomial invariants when Z is closed.

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The relative symplectic invariant of (Z, R, F) depends only on F through its homotopy class.

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- Define \mathcal{N} for fiber bundles.

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- Define \mathcal{N} for fiber bundles.

This uses the Narasimhan-Seshadri correspondence for surfaces.

- Discuss non-fiber bundle case.

This is where quilts enter the picture (e.g., the *quilted Atiyah-Floer conjecture*, and its 4-manifold generalizations).

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$P \rightarrow \Sigma$ as above

$$\mathcal{A}(\Sigma) = \{\text{connections on } P\}$$

$$\mathcal{A}_{\text{flat}}(\Sigma) = \{\alpha \in \mathcal{A}(\Sigma) \mid F_\alpha = 0\}$$

where F_α is the curvature of α

$$\mathcal{G}(\Sigma) = \{\text{gauge transformations on } P\}$$

$$\mathcal{G}_0(\Sigma) = \{g \in \mathcal{G}(\Sigma) \mid g \simeq \text{Id}\}$$

The action of $\mathcal{G}_0(\Sigma)$ on $\mathcal{A}_{\text{flat}}(\Sigma)$ is free, and

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Complexified gauge group

Recall $\mathcal{G}(\Sigma) = \{\text{sections of } P \times_G G \rightarrow \Sigma\}$

$G^{\mathbb{C}}$ complexification of $G = \text{SO}(3)$

$\mathcal{G}^{\mathbb{C}} = \{\text{sections of } P \times_G G^{\mathbb{C}} \rightarrow \Sigma\}$

$\mathcal{G}_0^{\mathbb{C}} := \{g \in \mathcal{G}^{\mathbb{C}} \mid g \simeq \text{Id}\}$

These act on $\mathcal{A}(\Sigma)$ by identifying $\mathcal{A}(\Sigma)$ with the space of Cauchy-Riemann operators.

There is a preferred neighborhood $\mathcal{A}^{ss}(\Sigma)$ of $\mathcal{A}_{\text{flat}}(\Sigma)$ on which $\mathcal{G}_0^{\mathbb{C}}$ acts freely.

There is some $\epsilon_0 > 0$ so that if $\|F_{\alpha}\|_{L^2} \leq \epsilon_0$, then $\alpha \in \mathcal{A}^{ss}(\Sigma)$.

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These act on $\mathcal{A}(\Sigma)$ by identifying $\mathcal{A}(\Sigma)$ with the space of Cauchy-Riemann operators.

There is a preferred neighborhood $\mathcal{A}^{ss}(\Sigma)$ of $\mathcal{A}_{\text{flat}}(\Sigma)$ on which $\mathcal{G}_0^{\mathbb{C}}$ acts freely.

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Complexified gauge group

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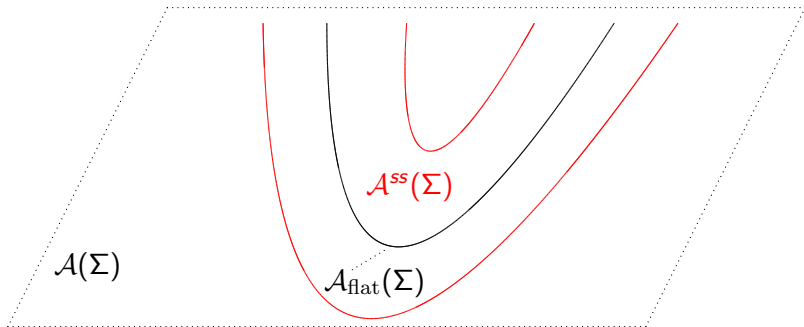
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$$\frac{\mathcal{A}^{SS}(\Sigma)}{\mathcal{G}_0^C}$$

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Theorem (Narasimhan-Seshadri (1965))

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$$M(\Sigma) := \frac{\mathcal{A}_{\text{flat}}(\Sigma)}{\mathcal{G}_0(\Sigma)} \cong \frac{\mathcal{A}^{SS}(\Sigma)}{\mathcal{G}_0^{\mathbb{C}}}$$

Denote the projection by

$$\Pi : \mathcal{A}^{SS}(\Sigma) \longrightarrow M(\Sigma)$$

A property of Π

- The metric and orientation on Σ induce a Hodge star $*$

- $*^2 = -1$ on 1-forms

- $T_\alpha \mathcal{A}(\Sigma) = \{\text{Lie algebra-valued 1-forms}\}$

$\Rightarrow \mathcal{A}(\Sigma)$ has a complex structure given by $*$

Similarly, $M(\Sigma)$ has a complex structure J induced from $*$.

Π preserves this structure:

$$D_\alpha \Pi(*\mu) = J D_\alpha \Pi(\mu)$$

where $D_\alpha \Pi$ is the linearization at $\alpha \in \mathcal{A}^{ss}(\Sigma)$ and $\mu \in T_\alpha \mathcal{A}(\Sigma)$.

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- The map Π takes holomorphic curves in $\mathcal{A}^{SS}(\Sigma)$ to holomorphic curves in $M(\Sigma)$.

- This can be strengthened as follows:

The tangent space to the complex gauge orbit through α is

$$\ker D_\alpha \Pi = \text{Im } d_\alpha \oplus \text{Im } * d_\alpha,$$

where d_α is the covariant derivative of α in the adjoint representation.

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In particular, if $\alpha : S \subseteq \mathbb{C} \rightarrow \mathcal{A}(\Sigma)$ satisfies

$$\partial_S \alpha + * \partial_t \alpha \in \text{Im } d_\alpha \oplus \text{Im } * d_\alpha, \quad \text{and}$$

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Σ as above

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$Z := S \times \Sigma$ with induced orientation and metric

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Local picture: The case $S \times \Sigma$

Suppose $A \in \mathcal{A}(S \times \Sigma)$ is a connection.

(s, t) coordinates on $S \subset \mathbb{C}$.

Write A in coordinates

$$A = \alpha + \phi ds + \psi dt,$$

$\alpha : S \rightarrow \mathcal{A}(\Sigma)$ defined by $\alpha(s, t) := A|_{\{(s,t)\} \times \Sigma}$

ϕ, ψ defined by contracting A with ∂_s, ∂_t , respectively.

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A connection A on Z is an **(ASD) instanton** if

$$*_Z F_A = -F_A,$$

where $*_Z$ is the Hodge star on Z .

$A = \alpha + \phi ds + \psi dt$ is an instanton if and only if

$$\begin{aligned}\partial_s \alpha + *_\Sigma \partial_t \alpha &= d_\alpha \phi + *_\Sigma d_\alpha \psi \\ F_\alpha &= *_\Sigma (\partial_t \phi - \partial_s \psi - [\phi, \psi])\end{aligned}$$

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In particular, if we can ensure the curvature F_α is small, then

$$u := \Pi(\alpha) : S \longrightarrow M(\Sigma)$$

is a well-defined function that is holomorphic.

We are starting to see a map

$$\{\text{instantons}\} \longrightarrow \{\text{holomorphic curves } S \rightarrow M(\Sigma)\},$$

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Now assume we have a fiber bundle

$$\Sigma \hookrightarrow Z \xrightarrow{F} S.$$

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Global picture

There exist perturbations that yield smooth moduli spaces (of instantons and holo. curves).

Assume everything is suitably perturbed.

The instanton moduli space

$$\{A \in \mathcal{A}(Z) \mid *F_A = -F_A\} / \mathcal{G}_0$$

is a smooth manifold.

We will denote the dimension near $[A]$ by $\dim(A)$.

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Global picture

Theorem (D. (2013))

Suppose S is closed or has cylindrical ends. Given any $\epsilon_0 > 0$, there is a metric on S so that if A is an instanton on Z with $\dim(A) \leq 3$, then

$$\sup_{x \in S} \|F_\alpha\|_{L^2(\Sigma_x)} \leq \epsilon_0.$$

That is, $\alpha(x) = A|_{\Sigma_x}$ is in the domain of Π for all $x \in S$.

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Define

$$\mathcal{N} : \{\text{instantons } A \text{ on } Z \text{ with } \dim(A) \leq 3\} / \{\text{gauge}\} \\ \longrightarrow \{\text{holomorphic sections } S \rightarrow M(Z)\}.$$

by sending $[A]$ to the map $x \mapsto \Pi(\alpha(x))$.

Remarks

- The map \mathcal{N} preserves the dimension of the moduli spaces.
- When S has cylindrical ends then \mathcal{N} preserves the broken trajectory compactifications of the moduli spaces.

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Theorem (D.-McNamara (2014))

The map \mathcal{N} is a C^1 -embedding.

Proof.

Injectivity of \mathcal{N} : If $\mathcal{N}([A_0]) = \mathcal{N}([A_1])$, then $\Pi(\alpha_0) = \Pi(\alpha_1)$ as maps $S \rightarrow M(Z)$. This implies

$$\alpha_0 = g^* \alpha_1$$

for some complex gauge transformation $g : S \rightarrow \mathcal{G}_0^{\mathbb{C}}$. The instanton equation and irreducibility force g to take values in the real gauge group \mathcal{G}_0 . This implies A_0 and A_1 lie in the same gauge orbit, so $[A_0] = [A_1]$.

Injectivity of the linearization $D_{[A]} \mathcal{N}$ is similar.



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Corollary (Dostoglou-Salamon, D.-McNamara, D.-)

The map \mathcal{N} is a diffeomorphism.

Proof 1 (Kähler case).

If Z is Kähler, then there is a complexified gauge group $\mathcal{G}(Z)^{\mathbb{C}}$. Any holomorphic section can be represented by a connection A_0 on Z that is in the semistable range. Then there is a (unique up to real gauge transformation) instanton A that is in the $\mathcal{G}(Z)^{\mathbb{C}}$ -orbit of A_0 . The map sending $[A_0]$ to $[A]$ is an approximate inverse of \mathcal{N} . □

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Proof 2 (general case).

Given a connection A_0 that represents a holomorphic section of $M(Z)$, one can use an implicit function theorem as in Dostoglou-Salamon [DS3] to find a nearby instanton A . This provides an approximate inverse to \mathcal{N} . \square

Proof 3 (general case).

Given a connection A_0 that represents a holomorphic section of $M(Z)$, one can use the Yang-Mills heat flow to flow down to an instanton A . This is another approximate inverse of \mathcal{N} . \square

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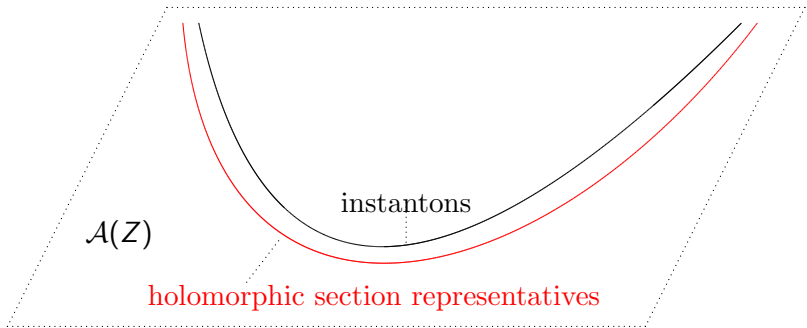
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Set-up (reminder)

Suppose

- S is a connected, oriented surface with cylindrical ends,
- Z is an oriented 4-manifold with the same number of cylindrical ends as S ,
- $F : Z \rightarrow S$ preserves the ends and has connected, nonempty fibers, and
- $R \rightarrow Z$ is a principal $\mathrm{SO}(3)$ -bundle that restricts to the non-trivial bundle on some fiber.

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Compare instantons on Z with certain symplectic objects, and try to extend the results above.

Instantons are well-defined, and the (relative) Donaldson invariant depends only on Z , the topological type of the bundle R , and the homotopy class of F .

The relevant symplectic objects are holomorphic *quilts*. These were defined by Wehrheim-Woodward [WW2], and one uses *Morse 2-functions* to obtain these from 4-manifolds.

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Morse 2-functions

$F : Z \rightarrow S$ is a **Morse 2-function** if it can be written locally as a ‘homotopy of Morse functions’. That is, if each point in S lies in a neighborhood $I \times I \subset S$, with $I = [0, 1]$, and this neighborhood satisfies

- $F^{-1}(I \times I) = I \times Y$ where Y is a 3-manifold,
- in these coordinates

$$F(s, y) = (s, f_s(y)),$$

where $f_s : Y \rightarrow I$ is Morse for all but a finite set of exceptional times s , and

- each exceptional time corresponds to a critical point birth, death or crossing.

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Morse 2-functions

Example 1: Suppose $F : Y \rightarrow S$ is a fiber bundle.

Example 2: Suppose Y is a closed 3-manifold and $f : Y \rightarrow S^1$ is a Morse function. Consider $Z = \mathbb{R} \times Y$ and $S = \mathbb{R} \times S^1$. Then $F(s, y) = (s, f(y))$ is a Morse 2-function. This is the **cylindrical case**.

In general, if Z has cylindrical ends, then we assume that any Morse 2-function restricts to the cylindrical case on the ends.

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It is useful to trace out the critical values of the f_S in the surface S . This is called the **Cerf graphic**. The critical values form the **seams**, and the regular values form the **patches**.

Morse 2-functions

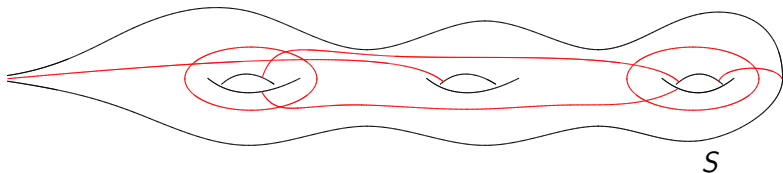
In the cylindrical case $Z = \mathbb{R} \times Y$, the Cerf graphic takes the form:



$$S = \mathbb{R} \times S^1$$

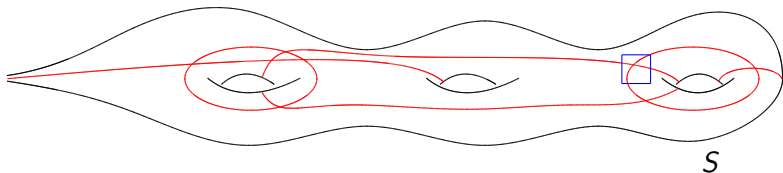
Morse 2-functions

Here is the Cerf graphic for some Morse 2-function $F : Z \rightarrow S$, where S is a punctured genus 3 surface.



Morse 2-functions

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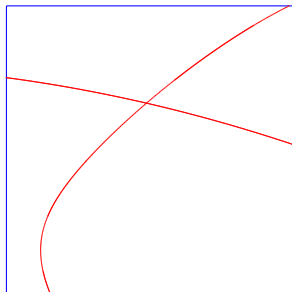
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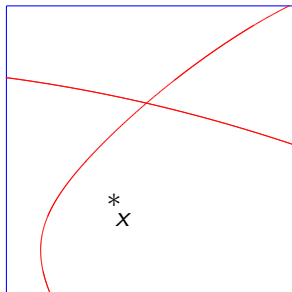
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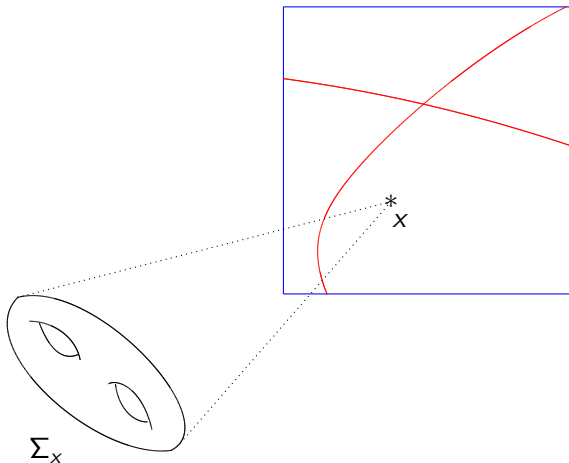
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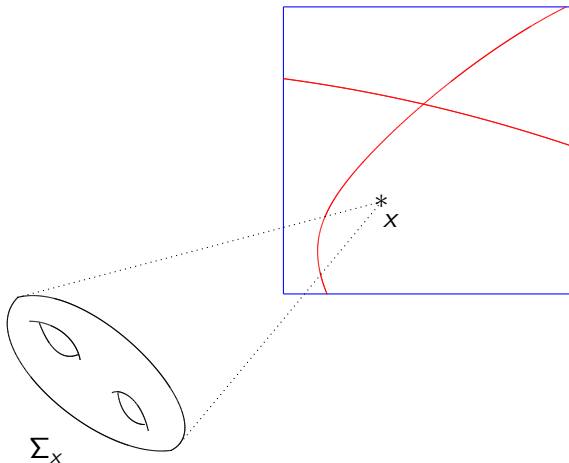
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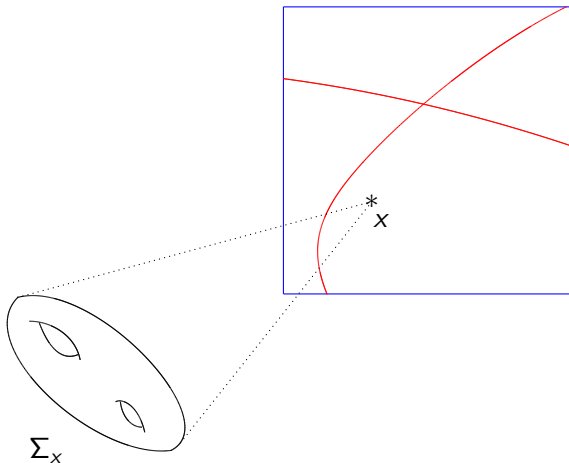
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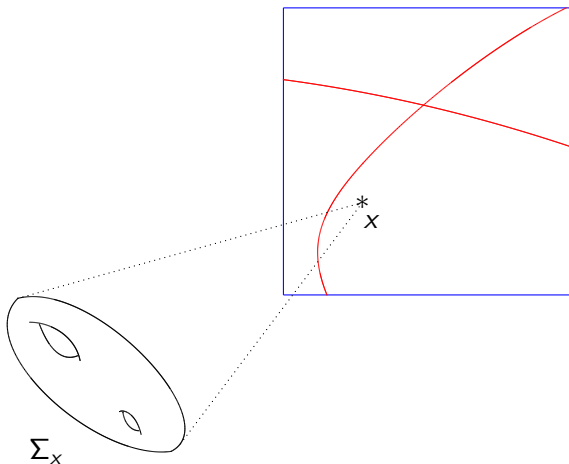
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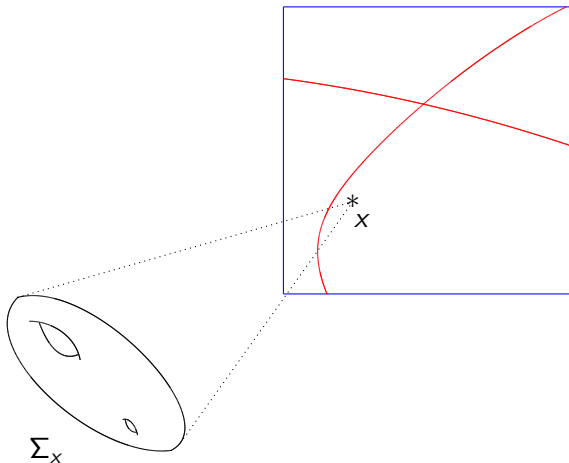
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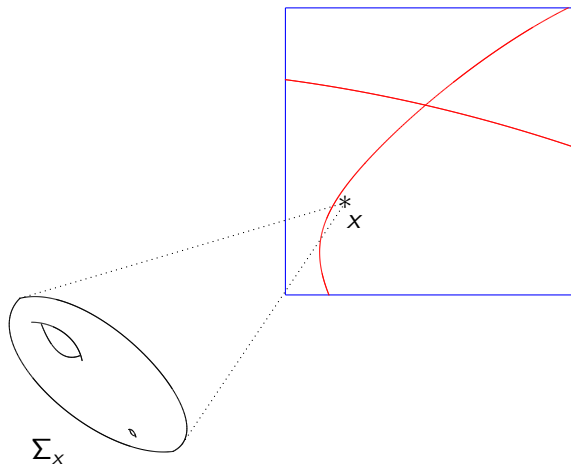
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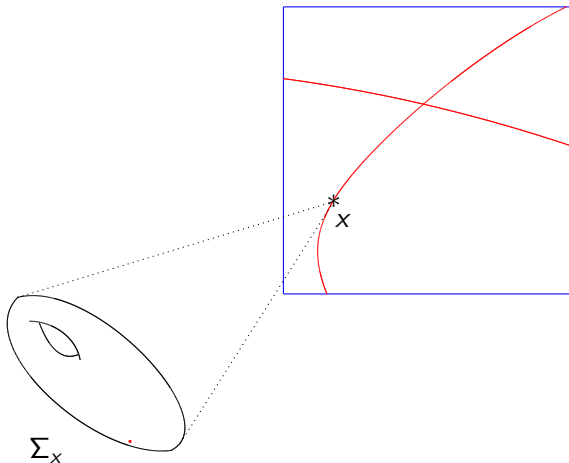
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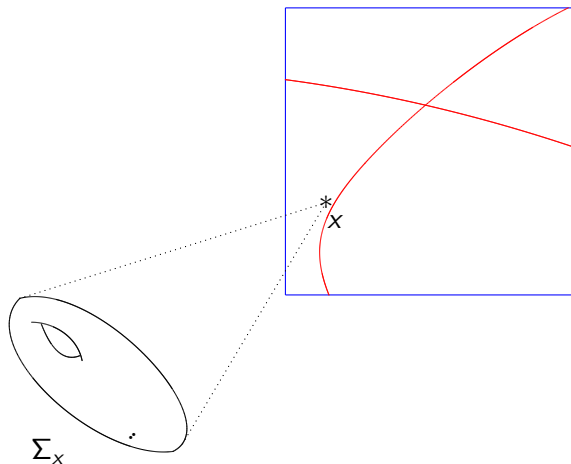
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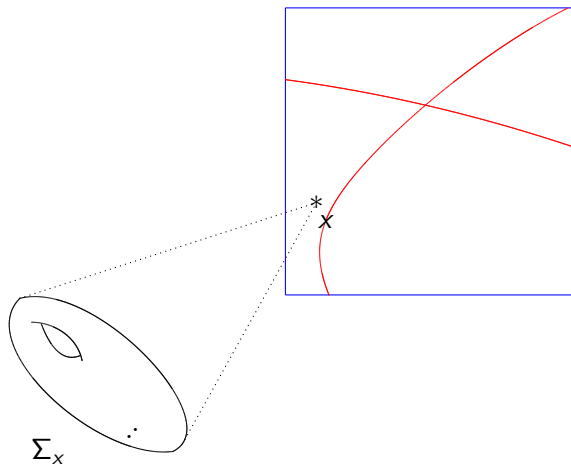
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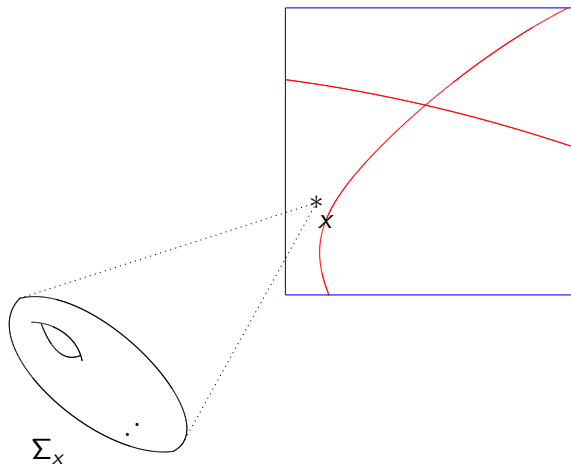
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Morse 2-functions

Theorem (Gay-Kirby (2011))

Suppose $F : Z \rightarrow S$ is any smooth function with connected and non-empty fibers. Then F is homotopic to a Morse 2-function that has connected and non-empty fibers. This Morse 2-function is unique up to certain homotopies of Morse 2-functions.

Quilts from Morse 2-functions

The treatment here follows Wehrheim-Woodward [WW1] and Wehrheim [W].

Assume $F : Z \rightarrow S$ is a Morse 2-function.

Consider the surface S with the diagram coming from the critical points.

Each point x in a patch is a regular point for F . The inverse image $F^{-1}(x) = \Sigma_x$ is therefore a closed, connected, oriented surface. Label this point with the symplectic manifold $M(\Sigma_x)$.

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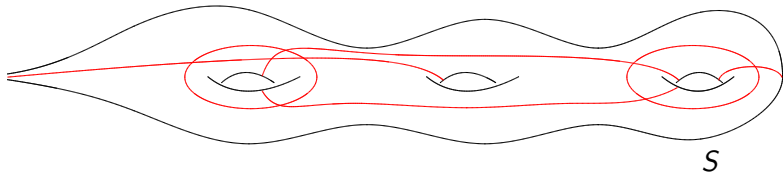
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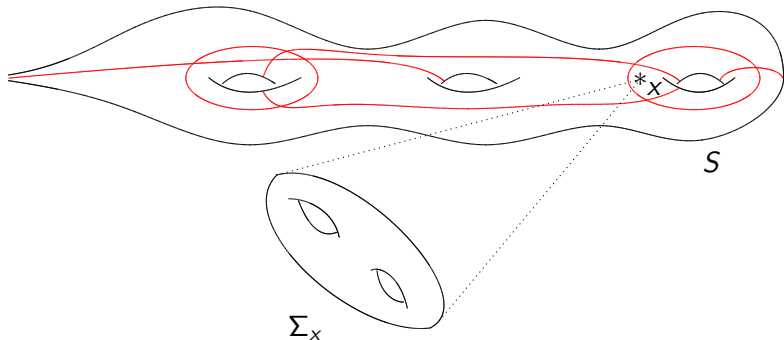
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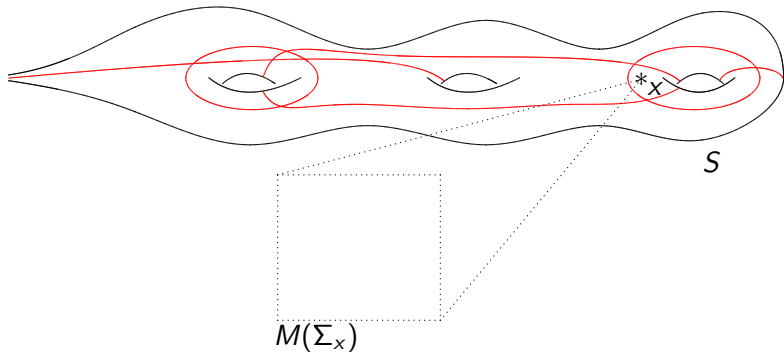
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Quilts from Morse 2-functions

Over each patch S_i is a fiber bundle, as before: $M(Z_i) \rightarrow S_i$

Think of the singular fibers as interpolating between non-singular fibers. (In symplectic language, there is a *Lagrangian correspondence* here.)

A **quilt** is a collection $\{u_i\}$ where u_i is a section $S_i \rightarrow M(Z_i)$, and we require that these agree on the seams.

The quilt is **holomorphic** if each u_i is holomorphic with respect to some complex structure.

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Wehrheim-Woodward show the following:

There is a complex structure on each $M(Z_i)$ such that, after perturbing, the moduli space of holomorphic quilts is smooth and finite-dimensional.

Counting the elements of the zero-dimensional component of this moduli space defines the relevant symplectic invariant.

In the cylindrical case, this invariant depends only on $\mathbb{R} \times Y$, the isomorphism class of R , and the homotopy class of f .

(For more general 4-manifolds, it is not yet known whether this depends only on the homotopy class of F . Establishing this is an active research project of Wehrheim.)

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Wehrheim-Woodward show the following:

There is a complex structure on each $M(Z_i)$ such that, after perturbing, the moduli space of holomorphic quilts is smooth and finite-dimensional.

Counting the elements of the zero-dimensional component of this moduli space defines the relevant symplectic invariant.

In the cylindrical case, this invariant depends only on $\mathbb{R} \times Y$, the isomorphism class of R , and the homotopy class of f .

(For more general 4-manifolds, it is not yet known whether this depends only on the homotopy class of F . Establishing this is an active research project of Wehrheim.)

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Suppose A is an instanton. For each patch S_i and $x \in S_i$, set $\alpha_i(x) = A|_{\{F^{-1}(x)\}}$.

As before, there is a metric on Z for which each $\alpha_i(x)$ has small curvature.

Define $u_i(x) := \Pi(\alpha_i(x))$. So u_i is a section $S_i \rightarrow M(Z_i)$.

Then the fiber bundle argument from above shows that u_i is holomorphic in the interior of the patch S_i .

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That is, just as before, we should get a map from the moduli space of instantons to the moduli space of holomorphic quilts.

The catch is that the relevant metric on each fiber Σ_x becomes singular (blows up) as x approaches a seam. This reflects the fact that the fibers over the seams are singular.

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Q1: Do the u_i actually have the expected boundary conditions?

This should be the case. The work is in studying how the map $\Pi : \mathcal{A}^{SS}(\Sigma) \rightarrow M(\Sigma)$ behaves under a degeneration of the metric.

If this is the case, then the relevant holomorphic quilts would be holomorphic with respect to a complex structure that *becomes singular near the boundary*.

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Q2: Does the associated moduli space have the regularity/compactness properties needed to obtain a well-defined invariant?

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Supposing these questions have favorable answers, then we should be able to extend the analysis from the fiber bundle picture above to show that the induced map \mathcal{N} is a diffeomorphism from the moduli space of instantons to the moduli space of holomorphic quilts.

Even if these questions do not have favorable answers, there is already enough going on here to obtain strong results. For example, these ideas were used to prove a compactness theorem [D1] that studies the behavior of instantons under certain degenerations of the underlying metric.

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Thank you for your attention.

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