

Adjoint based parameter recovery

*or: the beauty of
Integration by Parts*

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Acknowledgments

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Outline

- Introduction
- Why?
- Adjoint Theory
- Some numerics
- Conclusions

Apology

We'd like to study the Richards Equation:

$$C(h)\partial_t h = \partial_x(K(h)(\partial_x h - 1))$$

where $C(h)$ is soil capacity and $K(h)$ is hydraulic conductivity. But...

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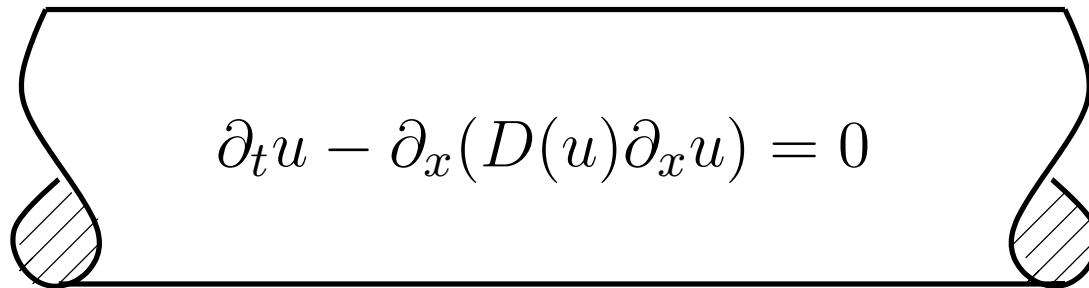
we'll consider the simpler problem

$$\partial_t u = \partial_x(D(u)\partial_x u),$$

which captures many details of the method.

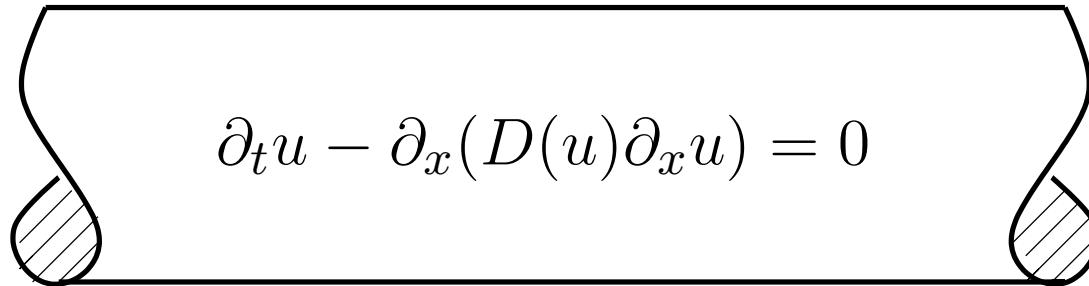
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The physical set-up



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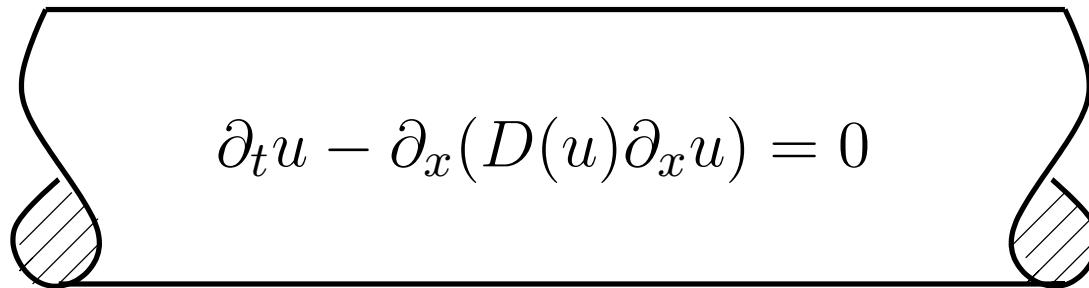


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$$\partial_x u(1, t) = 0$$

Introduction

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GOAL:

Given some output measurements of this system, recover $D(u)$.

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Let W represent this class of functions.

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- Why the toy problem?
- Why adjoint methods?

Forward

The (Forward) Problem: Given $D(u)$ and $f(t)$, find u which satisfies:

$$\begin{aligned}\partial_t u - \partial_x(D(u)\partial_x u) &= 0 & (x, t) \in U_T \\ u(0, t) &= f(t) & t \in (0, T) \\ \partial_x u(1, t) &= 0 & t \in (0, T) \\ u(x, 0) &= f(0) & x \in (0, 1)\end{aligned}$$

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Why?

- Unique weak (forward) solution if D and f from W .
- Model of a simple physical process.
- Measurements of flux at $x = 0$ and state at $x = 1$ possible.

Inverse

The (Inverse) Problem: Given boundary measurements $g(t)$ and $h(t)$, recover $D(u)$ which satisfies:

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Why?

- Experimental interest.
- Simple measurements.
- If we can recover a D from W , it should be unique.

Adjoint

Max Min Statement

- for each t , $f(0) < u(x, t) < f(t)$, $0 < x < 1$
- $\partial_x u(x, t) < 0$ a.e. on $(0, 1) \times (0, T)$

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Comments:

If we let $B(u) = \int_0^u D(s) \, ds$, then $D(u)\partial_x u = \partial_x B(u)$.

Also, $B(u) = \int_0^u D(s) \, ds = (u - 0) D(\tilde{u})$ for some \tilde{u} between u and 0.

Adjoint

Again, let $g(t) := -D(u(0, t))\partial_x u(0, t)$ and $\Phi[D, f] = g(t)$.

We can show that if $D_1(u) > D_2(u)$ then the map Φ is:

- Monotone : $\Phi[D_1, f](t) > \Phi[D_2, f](t)$

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- Continuous: $\|g_1 - g_2\|_{L^2(0,T)} \leq C\|D_1 - D_2\|_\infty$
- Injective : $\Phi[D_1, f] = \Phi[D_2, f]$ implies $D_1 = D_2$

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Adjoint

Integral Identity:

Let u, v represent solutions to direct problem with $D = D_1, D_2$, respectively. Subtract, multiply by smooth ϕ and integrate by parts to yield:

$$\begin{aligned} & \iint_{U_T} (u - v)(\phi_t + k\phi_{xx}) - \int_0^1 (u - v)\phi \Big|_{t=0}^{t=T} + \\ & + \int_0^T (D_1(u)u_x - D_2(v)v_x)\phi \Big|_{x=0}^{x=1} + \int_0^T (D_1(u) - D_1(v))(k\phi_x) \Big|_{x=0}^{x=1} \\ & = \iint_{U_T} (D_1(v) - D_2(v)v_x)\phi_x \end{aligned}$$

Here $U_T = (0, 1) \times (0, T)$ and $k := D(\tilde{u})$ for some \tilde{u} in the interval (v, u)

Adjoint

The general adjoint problem:

$$\begin{aligned}\partial_t \phi + k \partial_{xx} \phi &= F^*(x, t) & (x, t) \in U_T \\ \phi(x, T) &= p^*(x) & x \in (0, 1) \\ k \phi_x &= g^*(t) & x = 0, t \in (0, T). \\ \phi &= h^*(t) & x = 1, t \in (0, T)\end{aligned}$$

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Adjoint

The dual data acts like switches. Data appears in a duality pairing with terms of the direct problem.

Taking $F^* = 0, p^* = 0, h^* = 0$, the identity collapses to

$$\int_0^T g^*(g_1 - g_2) = \iint_{U_T} (D_1(v) - D_2(v)) v_x \phi_x$$

Adjoint

Letting $\Phi[D, f] = g(t)$, we could cast this formally as:

$$\begin{aligned} (\Phi[D_1, f] - \Phi[D_2, f], \theta)_{L^2} &:= (\delta\Phi[D_1, D_2]\Delta D, \theta)_{L^2} \\ &= \langle \Delta D, {}^t\delta\Phi[D_1, D_2] \theta \rangle_{W \times W^*} \end{aligned}$$

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The integral identity provides a realization of this expression.

$$\int_0^T g^*(g_1 - g_2) = \iint_{U_T} (D_1(v) - D_2(v)) v_x \phi_x$$

Adjoint

Now some quick proofs:

$$\int_0^T g^*(g_1 - g_2) dt = \iint_{U_T} (D_1(v) - D_2(v)) v_x \phi_x dx dt$$

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- **Continuous:** Take $g^* = (g_1 - g_2)/\|g_1 - g_2\|_{L^2}$.

Adjoint

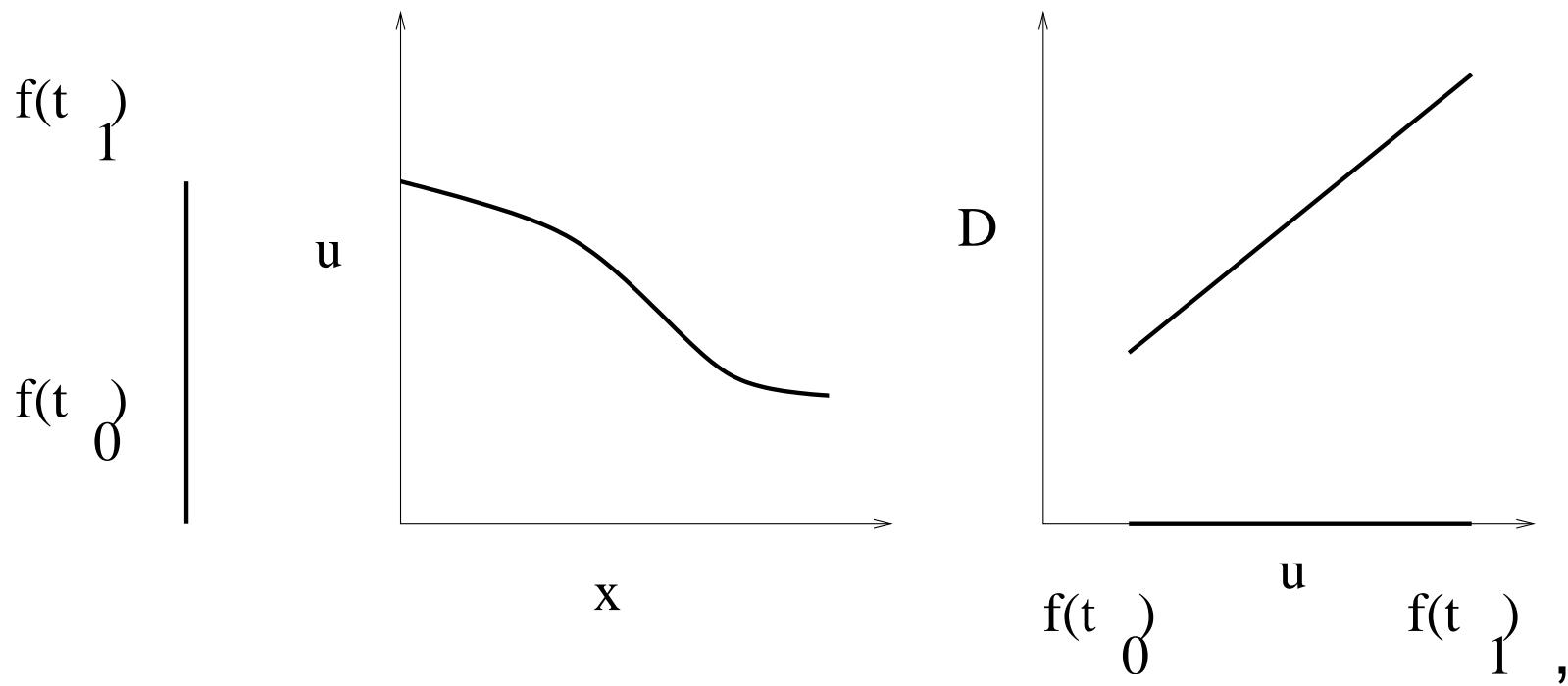
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- **Continuous:** Take $g^* = (g_1 - g_2)/\|g_1 - g_2\|_{L^2}$.
- **Injective:** Assume $D_1 - D_2 \not\equiv 0$. This implies that $g_1 - g_2 \not\equiv 0$.

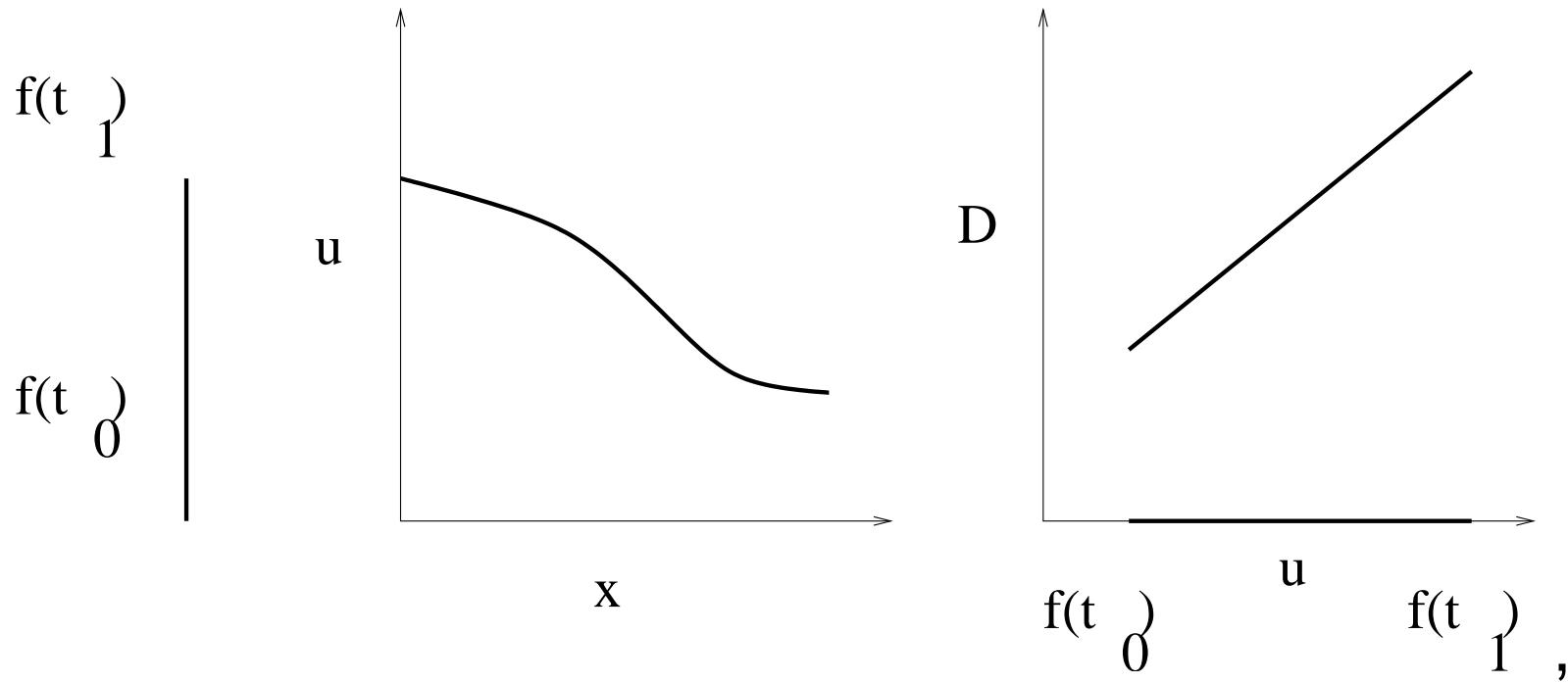
Numerics

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and so $[t_0, t_1] \Rightarrow [u_0, u_1]$. If

$$\hat{D}(u) = \sum_{k=0}^N d_k \Lambda_k,$$

Numerics

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$$(d_1 - \delta_1) = \frac{\int_0^T g^*(g_1 - g_2) dt}{\iint_{U_t} \Lambda u_x \phi_x dx dt}$$

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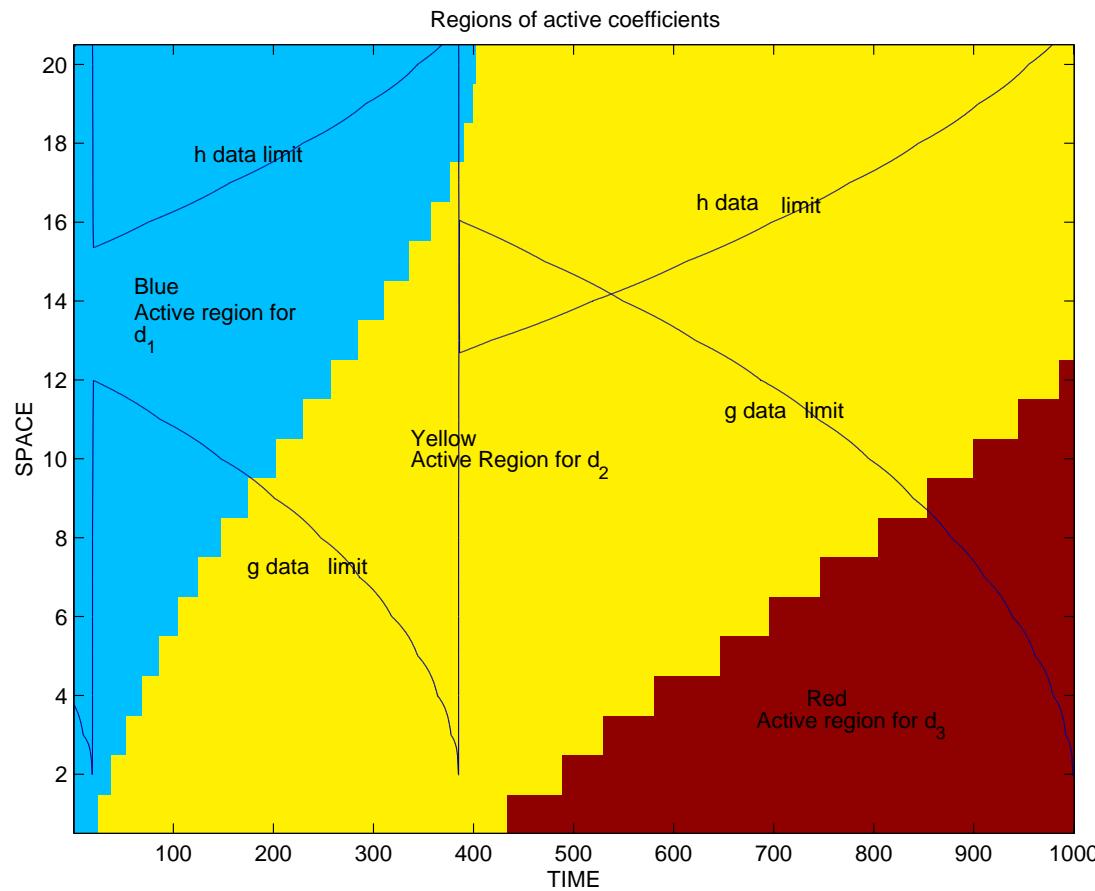
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- Apply integral ID to compute update for D .
- Repeat

Numerics

Regions of coefficient activity



Numerics

Apply identity to region where only d_1 is active.

$$b_1 = \int_0^{t_1} g^* \Delta g = (d_1 - \delta_1) \iint_{R_{11}} \Lambda_1(u) u_x \phi_x = A_{11} \Delta d_1$$

where $R_{11} = \{(x, t) \in U_{t_1} : u_0 < u(x, t) < u_1\}$

Numerics

Apply identity to the U_{t_2} time strip. Only d_1 and d_2 are active.

$$\begin{aligned} b_2 := \int_{t_1}^{t_2} g^* \Delta g &= (d_2 - \delta_2) \iint_{R_{22}} \Lambda_2(u) u_x \phi_x + \\ &\quad + (d_1 - \delta_1) \iint_{R_{21}} \Lambda_1(u) u_x \phi_x \\ &= A_{22} \Delta d_2 + A_{21} \Delta d_1 \end{aligned}$$

with

$$\begin{aligned} R_{21} &= \{(x, t) \in U_{t_2} : u_0 < u < u_1\} \\ R_{22} &= \{(x, t) \in U_{t_2} : u_1 < u < u_2\} \end{aligned}$$

Numerics

By the Q_{t_N} strip, a diagonal system forms.

$$\begin{bmatrix} A_{11} & 0 & \cdots & 0 \\ A_{21} & A_{22} & \cdots & 0 \\ \vdots & \vdots & \ddots & 0 \\ A_{N1} & \cdots & \cdots & A_{NN} \end{bmatrix} \begin{bmatrix} \Delta d_1 \\ \Delta d_2 \\ \vdots \\ \Delta d_N \end{bmatrix} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_N \end{bmatrix}$$

where $A_{ij} = \iint_{R_{ij}} \Lambda_j(u) u_x \phi_x$ and $b_i = \int_{t_{i-1}}^{t_i} g^* \Delta g$

Numerics

Now for some demos:

- Is some data better?
- Noise?
- ????

Conclusions

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