# Stability Analysis via Hill's method or: Can we find stable wave forms? 

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## Introduction

We'll talk about computing the spectra of linear operators, including the associated eigenfunctions.

Why is this important?
Spectral Stability: Given an equilibrium solution, is it stable under perturbation? If so, it is a candidate for a solution of permanent form.

## Patterns in waves



February, 2003 at Maalea Bay on Maui. Courtesy of Robert I. Odom, Applied Physics Laboratory, University of Washinton

## More patterns



CERC harbor model (Source unknown)

## Spectral Stability

Consider the evolution system

$$
u_{t}=N(u)
$$

with an equilibrium solution $u_{e}$ :

$$
N\left(u_{e}\right)=0 .
$$

Is this solution stable or unstable?
Linear analysis: let

$$
u=u_{e}+\epsilon \psi .
$$

Substitute in and retain first-order terms in $\epsilon$ :

$$
\psi_{t}=\mathcal{L}\left[u_{e}(x)\right] \psi .
$$

## Eigenfunction expansion

Separation of variables: $\psi(x, t)=e^{\lambda t} z(x)$ :

$$
\mathcal{L}\left[u_{e}(x)\right] z=\lambda z .
$$

- This is a spectral problem.
- If $\Re(\lambda) \leq 0$ for all bounded $z(x)$, then $u_{e}$ is spectrally stable.


## Application

Our starting point is

$$
\mathcal{L} z=\lambda z,
$$

with

$$
\mathcal{L}=\sum_{k=0}^{M} f_{k}(x) \partial_{x}^{k}, \quad f_{k}(x+L)=f_{k}(x) .
$$

We want to find

- Spectrum $\sigma(\mathcal{L})=\left\{\lambda \in \mathbb{C}:\|z\|_{\infty}<\infty\right\}$.
- Corresponding eigenfunctions $z(\lambda, x)$ ?


## Floquet's Theorem

Consider

$$
\begin{equation*}
\varphi_{x}=A(x) \varphi, \quad A(x+L)=A(x) . \tag{*}
\end{equation*}
$$

Floquet's theorem states that the fundamental matrix $\Phi$ for this system has the decomposition

$$
\Phi(x)=P(x) e^{R x},
$$

with $P(x+L)=P(x)$ and $R$ constant.

## Floquet's Theorem

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$$
\Phi(x)=P(x) e^{R x},
$$

with $P(x+L)=P(x)$ and $R$ constant.
Conclusion: All bounded solutions of (*) are of the form

$$
\varphi=e^{i \mu x} \sum_{n=-\infty}^{\infty} \hat{\varphi}_{n} e^{i 2 \pi n x / L},
$$

with $\mu \in[0,2 \pi / L)$.

## Eigenfunctions

The periodic eigenfunctions can be expanded as

$$
\varphi=e^{i \mu x} \sum_{n=-\infty}^{\infty} \hat{\varphi}_{n} e^{i \pi n x / L}
$$

with $\mu \in[0, \pi / L)$
Substitute in the equation and cancel $e^{i \mu x}$.
The Floquet parameter $\mu$ only appears in derivative terms.

## Hill's method

- Find Fourier coefficients of all functions
- Choose a number of $\mu$ values $\mu_{1}, \mu_{2}, \ldots$
- For all chosen $\mu$ values, construct $\hat{\mathcal{L}}_{N}(\mu)$
- Use favorite eigenvalue/vector solver
- Reconstruct eigenfunctions corresponding to eigenvalues


## NLS

I've been looking at solutions of the 2-D cubic nonlinear Schrödinger (NLS) equation, given by

$$
i \phi_{t}+\alpha \phi_{x x}+\beta \phi_{y y}+\gamma|\phi|^{2} \phi=0 .
$$

The NLS equation arises in many models:

- Bose-Einstein condensates $(\alpha \beta>0)$
- Deep water models $(\alpha \beta<0)$
- Optics ( $\alpha \beta>0$ )


## NLS

Consider

$$
i \psi_{t}+\alpha \psi_{x x}+\beta \psi_{y y}+|\psi|^{2} \psi=0
$$

This equation has exact 1-D traveling wave solutions of the form

$$
\psi(x, t)=\phi(x) e^{i \omega t+i \theta(x)}
$$

where $\phi$ and $\theta$ are real-valued functions and $\omega$ is a real constant.

- If $\theta(x)=$ constant then the solution has trivial phase (TP).
- if $\theta(x) \neq$ constant the solution has nontrivial phase (NTP).


## NLS

More precisely,

$$
\psi(x, t)=\phi(x) e^{i \omega t+i \theta(x)}
$$

where

$$
\begin{aligned}
\phi^{2}(x) & =\alpha\left(-2 k^{2} \operatorname{sn}^{2}(x, k)+B\right), \\
\theta(x) & =c \int_{0}^{x} \phi^{-2}(\tau) d \tau, \\
\omega & =\frac{1}{2} \alpha\left(3 B-2\left(1+k^{2}\right)\right), \quad \text { and } \\
c^{2} & =-\frac{\alpha^{2}}{2} B\left(B-2 k^{2}\right)(B-2) .
\end{aligned}
$$

$k$ and $B$ are free parameters and $\mathrm{sn}(x, k)$ is the Jacobi elliptic sine function.

## Focusing

NLS is focusing or attractive in the $x$ dimension if $\alpha>0$. To make $\phi$ real in this case, we choose $B$ in $\left[2 k^{2}, 2\right]$.


## Defocusing

NLS is defocusing or repulsive if $\alpha<0$. To make $\phi$ real in this case, we choose $B \leq 0$.


## Linearized TP spectral problem

Now consider the (modulus and phase) perturbed TP solution of the form

$$
\psi_{p}=(\phi+\epsilon u+i \epsilon v) e^{i \omega t}
$$

Linearizing and considering real and imaginary contributions yields the system

$$
\begin{aligned}
\omega u-3 \phi^{2} u-\beta u_{y y}-\alpha u_{x x} & =-v_{t} \\
\omega v-\phi^{2} v-\beta v_{y y}-\alpha v_{x x} & =u_{t}
\end{aligned}
$$

## Linearized TP spectral problem

Let $u(x, y, t)=U(x) e^{i \rho y+\lambda t}$
and $v(x, y, t)=V(x) e^{i \rho y+\lambda t}$.
Then

$$
\begin{aligned}
\omega u-3 \phi^{2} u-\beta u_{y y}-\alpha u_{x x} & =-v_{t} \\
\omega v-\phi^{2} v-\beta v_{y y}-\alpha v_{x x} & =u_{t}
\end{aligned}
$$

becomes

$$
\begin{aligned}
\omega U-3 \phi^{2} U+\beta \rho^{2} U-\alpha U_{x x} & =-\lambda V \\
\omega V-\phi^{2} V+\beta \rho^{2} V-\alpha V_{x x} & =\lambda U
\end{aligned}
$$

## Linearized TP spectral problem

We write

$$
\begin{aligned}
\omega U-3 \phi^{2} U+\beta \rho^{2} U-\alpha U_{x x} & =-\lambda V \\
\omega V-\phi^{2} V+\beta \rho^{2} V-\alpha V_{x x} & =\lambda U
\end{aligned}
$$

as

$$
\mathcal{L}\left[\begin{array}{l}
U \\
V
\end{array}\right]:=\left[\begin{array}{cc}
0 & L_{-} \\
-L_{+} & 0
\end{array}\right]\left[\begin{array}{l}
U \\
V
\end{array}\right]=\lambda\left[\begin{array}{l}
U \\
V
\end{array}\right]
$$

where

$$
L_{+}=\omega-3 \phi^{2}+\beta \rho^{2}-\partial_{x x}
$$

and

$$
L_{-}=\omega-\phi^{2}+\beta \rho^{2}-\partial_{x x}
$$

This is our linearized TP spectral problem. The coefficients are periodic.

## Spectral Stability

We can finally consider the spectral stability of the periodic coefficient linear problem

$$
\mathcal{L}\left[\begin{array}{l}
U \\
V
\end{array}\right]:=\left[\begin{array}{cc}
0 & L_{-} \\
-L_{+} & 0
\end{array}\right]\left[\begin{array}{l}
U \\
V
\end{array}\right]=\lambda\left[\begin{array}{l}
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$$

where

$$
L_{+}=\omega-3 \phi^{2}+\beta \rho^{2}-\partial_{x x}
$$

and

$$
L_{-}=\omega-\phi^{2}+\beta \rho^{2}-\partial_{x x}
$$

We now build the matrix $\hat{L}_{-}(\mu)$. The same method generates $\hat{L}_{+}(\mu)$.

## SN plus

In the literature, you might find graphs for spectra associated to periodically perturbed TP solutions.


## SN plus

## But now we can compute "all" unstable modes.



## Linearized NTP spectral problem

Now consider the (modulus and phase) perturbed NTP solution of the form

$$
\psi_{p}=(\phi+\epsilon u+i \epsilon v) e^{i \lambda t+i \theta}
$$

As in the TP case, we linearize, separate variables and introduce transverse perturbation to generate the system

$$
\begin{gathered}
\left(\omega-3 \gamma \phi^{2}+\beta \rho^{2}\right) U \\
\left(\omega-\gamma \phi^{2}+\beta \rho^{2}\right) V
\end{gathered}
$$

$$
\begin{aligned}
& -\alpha U_{x x}=-\lambda V \\
& -\alpha V_{x x}=\lambda U
\end{aligned}
$$

What does this spectrum look like?

## Linearized NTP spectral problem

Now consider the (modulus and phase) perturbed NTP solution of the form

$$
\psi_{p}=(\phi+\epsilon u+i \epsilon v) e^{i \lambda t+i \theta}
$$

As in the TP case, we linearize, separate variables and introduce transverse perturbation to generate the system

$$
\begin{array}{r}
\left(\omega-3 \gamma \phi^{2}+\beta \rho^{2}\right) U+\alpha c^{2} \phi^{-4} U+\alpha c \phi_{x}^{-2} V+2 \alpha c \phi^{-2} V_{x}-\alpha U_{x x}=-\lambda V \\
\left(\omega-\gamma \phi^{2}+\beta \rho^{2}\right) V+\alpha c^{2} \phi^{-4} V-\alpha c \phi_{x}^{-2} U-2 \alpha c \phi^{-2} U_{x}-\alpha V_{x x}=\lambda U
\end{array}
$$

What does this spectrum look like?

## Linearized NTP spectral problem

There are a lot of parameters: for each pair $( \pm \alpha, \pm \beta)$, we can pick

- $k$ elliptic modulus)
- $B$ offset, (constrained by $k$ )
- $\rho \quad$ wavenumber of perturbation
- $\mu$ Floquet parameter


## Linearized NTP spectral problem

For $\alpha=-\beta=1$ the spectrum might look like:

for $k=0.5, B=1$ and $\rho=0$.

## Linearized NTP spectral problem

and the eigenfunction corresponding to dominate eigenvalue:

where red $=\Re(U)$, blue $=\Im(U)$, green $=\Re(V)$ and black $=$ $\Im(V)$.

## Linearized NTP spectral problem

$\alpha=\beta=1$

$(k, B) \in(0,1) \times\left(2 k^{2}, 2\right)$ for $\rho \in[0,3]$.

## Linearized NTP spectral problem

$\alpha=-\beta=1$

$(k, B) \in(0,1) \times\left(2 k^{2}, 2\right)$ for $\rho \in[0,3]$.

## Linearized NTP spectral problem

$-\alpha=\beta=1$

$(k, B) \in(0,1) \times[-0.5,0]$ for $\rho \in[0,3]$.

## Linearized NTP spectral problem

$$
-\alpha=-\beta=1
$$


$(k, B) \in(0,1) \times[-0.5,0]$ for $\rho \in[0,3]$.

## Conclusions

Samples suggest that:
All NTP solutions to the cubic NLS equation are unstable with respect to perturbation.

Hill's method:

- Allows non-periodic eigenfunctions
- Simple to implement


## Thanks!!!



